# Crystal Structure of $\mathrm{ErCd}_{3}$ and Its Isomorphous Compounds 

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#### Abstract

The crystal structure of $\mathrm{ErCd}_{3}$ has been solved using photographic data. The compound is orthorhombic, $a=7.094, b=10.678, c=4.804 \AA$, space group $\mathrm{Cmcm}, Z=4$. The packing arrangement is closely related to that of $\mathrm{GdCd}_{3}$, with the hexagonal $\mathrm{Ni}_{3} \mathrm{Sn}$ structure type. All the $\mathrm{MCd}_{3}$ compounds with $\mathrm{M}=\mathrm{Tb}, \mathrm{Dy}$, $\mathrm{Ho}, \mathrm{Tm}, \mathrm{Lu}$ and Y are isotypic with $\mathrm{ErCd}_{3}$.


During a study on the alloying behaviour of the rareearth metals with cadmium (Bruzzone, Fornasini \& Merlo, 1972), Tb, Dy, Ho, Er, Tm, Lu and Y were found to form isomorphous $\mathrm{MCd}_{3}$ phases. The purpose of the present work was the determination of their crystal structure.

The preparation of the compounds is described by Bruzzone et al. (1972). Several, single needle-shaped crystals of $\mathrm{ErCd}_{3}$ were examined with X-rays by Laue, rotating, Weissenberg and precession methods, using Cu or $\mathrm{Mo} K \alpha$ radiation. The intensity data were obtained from a well formed prismatic crystal, measuring $0.2 \times 0.02 \times 0.02 \mathrm{~mm}$. Integrated precession photographs, taken with filtered Mo radiation, provided 217 independent reflexions of which 129 observed reflexions were measured by a microdensitometer and 88 unobserved were assumed to have a value equal to half the least observable intensity. Lorentz and polarization corrections were applied, but as the crystal was approximately cylindrical and $\mu r \simeq 0.5$ for Mo radiation, no absorption correction was made.

All crystals showed orthorhombic symmetry and the lattice constants, subsequently refined from powder data, are:

$$
a=7.094, \quad b=10.678, \quad c=4.804 \AA .
$$

The conditions for possible reflexions are: $h k l$ present with $h+k=2 n$, and $h 0 l$ present with $l=2 n$, indicating the three possible space groups $C m c 2_{1}, C 2 c m$ and Cmem.

A pycnometric determination of the density gave a value $d_{\mathrm{exp}}=9.1 \mathrm{~g} . \mathrm{cm}^{-3}$ which agrees with the calculated value $d_{\text {calc }}=9.21 \mathrm{g.cm}{ }^{-3}$, based on four unit formulae per cell.

The comparison of the powder photographs of $\mathrm{ErCd}_{3}$ with those of the hexagonal $\mathrm{GdCd}{ }_{3}$ phase $\left(\mathrm{Ni}_{3} \mathrm{Sn}\right.$ structure type, $a=6.621, c=4.933 \AA ;$ Bruzzone, Fornasini \& Merlo, 1971), shows a certain resemblance. The cell constants of $\mathrm{ErCd}_{3}$ are related to $\mathrm{Ni}_{3} \mathrm{Sn}$-like $\operatorname{GdCd}_{3}$ as follows: $a \simeq a_{\text {hex }}, b \simeq a_{\text {hex }} / 3$ and $c \simeq c_{\text {hex }}$. Moreover, it was noticed that the intensities of the spots $h k l$ were equal to those of $h_{1}, l l+2$, apart from the Lp factor and normal decline, suggesting that the atoms in the cell lie on two planes perpendicular to [001] at a distance of $c / 2$. A trial structure in the Cmcm
space group, based on a simple orthorhombic distortion of the $\mathrm{Ni}_{3} \mathrm{Sn}$ type, gave a reasonable agreement between calculated and observed intensities. A modified version of the program ORFLS (Busing, Martin \& Levy, 1962), was used for the full-matrix least-squares refinement. The function minimized was $\sum w\left(\left|F_{o}\right|-\right.$ $\left.\left|K F_{c}\right|\right)^{2}$. In this calculation, amplitudes of 129 reflexions that were actually observed were given unit weight; those of reflexions too weak to be observed were given zero weight. The atomic scattering factors were corrected for the anomalous scattering with the values given in International Tables for X-ray Crystallography (1962). Eleven parameters were refined - four positional coordinates, six scale factors and an overall temperature factor. After a few cycles including all observed reflexions, an $R$ value of 0.067 was obtained, where $R=\Sigma| | F_{o}\left|-\left|F_{c}\right|\right| / \Sigma\left|F_{o}\right|$. At this stage it was observed that some strong reflexions, namely 002,004 , 006,221 and 040 , showed an observed structure factor lower than the calculated one. A secondary extinction correction was then applied, using the formula $\left|F_{\text {correct }}\right|=\left|F_{o}\right|\left(1+g I_{c}\right)$ where $I_{c}$ is the calculated intensity and $g$ is a parameter to be determined empirically. Refinement proceeded assuming isotropic temperature factors for each atom as variables and gave a final $R$ value of 0.060 . For all reflexions the $R$ value becomes $0 \cdot 105$. After the last cycle the shifts in the coordinates were nearly one per cent of their estimated standard deviations. In Table 1 are listed crystallographic data and parameters of $\mathrm{ErCd}_{3}$ with the corresponding standard deviations. Table 2 collects the final observed and calculated structure factors. No attempt was made to refine further by using anisotropic temperature factors as variables, nor was refinement attempted using the space groups $C m c 2_{1}$ and $C 2 \mathrm{~cm}$.

Table 1. Crystallographic data for $\mathrm{ErCd}_{3}$
The estimated standard deviations from the least-squares refinements are given in parentheses in units of the last significant figure of the parameter value.

$$
\text { Space group: } \mathrm{Cmcm} \text { ( No 63) }
$$ $a=7 \cdot 094, b=10 \cdot 678, c=4 \cdot 804 \AA$ $d_{\mathrm{exp}}=9 \cdot 1 \mathrm{~g} . \mathrm{cm}^{-3} ; d_{\mathrm{c}: \mathrm{LL}}=9 \cdot 21 \mathrm{~g} . \mathrm{cm}^{-3}$.

|  | Position | $x$ | $y$ | $z$ | $B\left(\AA^{2}\right)$ |
| :--- | :---: | :--- | :---: | :---: | :---: |
|  | 4 Er | $4(c)$ | 0 | $0.3679(4)$ | $\frac{1}{4}$ |
| $4.40(8)$ |  |  |  |  |  |
| $4 \mathrm{Cd}(1)$ | $4(c)$ | 0 | $0.8249(8)$ | $\frac{1}{4}$ | $0.75(13)$ |
| $8 \mathrm{Cd}(2)$ | $8(g)$ | $0.2129(6)$ | $0.0950(5)$ | $\frac{1}{4}$ | $0.71(10)$ |

Table 2. Observed and calculated structure factors of $\mathrm{ErCd}_{3}$

An asterisk indicates reflexions which were too weak to be measured.


The crystal-structure determination was carried out also for the isomorphous compound $\mathrm{YCd}_{3}$. Intensity data were collected from a crystal measuring $0.13 \times$ $0.03 \times 0.04 \mathrm{~mm}$ using the same procedure as for $\mathrm{ErCd}_{3}$. A least-squares refinement was made with only $53 \mathrm{ob}-$ served reflexions, starting with the positional parameters of $\mathrm{ErCd}_{3}$, single isotropic $B$ factors and one scale factor. An $R$ value of 0.066 was obtained at the end of the last cycle. In Table 3 are shown crystallographic data for $\mathrm{YCd}_{3}$.

Table 3. Crystallographic data for $\mathrm{YCd}_{3}$

|  | $\begin{gathered} \text { Space group: Cmcm (No. 63) } \\ a=7.044, b=10.864, c=4.837 \AA \\ d_{\text {exp }}=7.6 \mathrm{~g} . \mathrm{cm}^{-3} ; d_{\text {calc }}=7.64 \mathrm{~g} . \mathrm{cm}^{-3} . \end{gathered}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | Position | $x$ | $y$ | $z$ | $B\left(\AA^{2}\right)$ |
| 4 Y | 4(c) | 0 | 0.368 (3) | $\frac{1}{4}$ | 1.5 (9) |
| $4 \mathrm{Cd}(1)$ | 4(c) | 0 | $0 \cdot 822$ (2) | 4 | $0 \cdot 8$ (5) |
| $8 \mathrm{Cd}(2)$ | 8(g) | $0 \cdot 216$ (2) | 0.094 (2) | $\frac{1}{4}$ | $0 \cdot 6$ (3) |


$z=\frac{1}{4}$

$z=\frac{3}{4}$

Fig. 1. The two layers at $z=\frac{1}{4}$ and $z=\frac{3}{4}$ of the elementary cell of $\mathrm{ErCd}_{3}$. Shaded circles: Er ; open circles: Cd.


Fig. 2. Geometrical relation between the $\mathrm{GdCd}_{3}\left(\mathrm{Ni}_{3} \mathrm{Sn}\right.$-type) and $\mathrm{ErCd}_{3}$ structures. Section $x y$ at $z=\frac{1}{4}$. Large circles are rare earth atoms, small circles Cd atoms. Four $\mathrm{GdCd}_{3}$ cells are drawn with solid lines; the dashed lines limit the $\mathrm{ErCd}_{3}$ cell.


Fig. 3. Atomic distribution around the rare earth and cadmium atoms in $\mathrm{GdCd}_{3}$ and $\mathrm{ErCd}_{3}$.

Fig. 1 shows the two layers at $z=\frac{1}{4}$ and $z=\frac{3}{4}$ of the elementary cell of $\mathrm{ErCd}_{3}$. In Table 4 are reported the interatomic distances for $\mathrm{ErCd}_{3}$ and $\mathrm{YCd}_{3}$. The very close relationship between the $\mathrm{GdCd}_{3}$ and $\mathrm{ErCd}_{3}$ structures is shown in Fig. 2. The arrows indicate the shifting directions of the atoms going from the hexagonal to the orthorhombic packing. In the $\mathrm{GdCd}_{3}$ and $\mathrm{ErCd}_{3}$ structures the coordination around the rare earth and cadmium atoms is somewhat different. This can be shown by plotting, for each atom in the asym-
metric unit, the number of atoms at equal distance vs. $d / \sum r$, where $d$ is the corresponding distance and $\Sigma r$ is the sum of the metallic radii of the considered atom and the surrounding one. As already noted by Bruzzone, Fornasini \& Merlo (1970), in all cases a gap in this atomic distribution occurs. The coordination numbers obtained by counting all the atoms before the gap are generally in good agreement with the values obtained following the Frank \& Kasper (1958, 1959) criterion. Fig. 3 shows the atomic distribution plot for $\mathrm{GdCd}_{3}$ and $\mathrm{ErCd}_{3}$. In $\mathrm{GdCd}_{3}$ the Gd atom is surrounded by twelve Cd atoms, six at the same distance and six at a different distance. In $\mathrm{ErCd}_{3}$ the coordination number of the rare-earth atom increases to 14 because an Er atom is bound with 12 Cd atoms and with two Er atoms. Regarding the two crystallographic types of cadmium in $\mathrm{ErCd}_{3}$, their coordination numbers are 12 and 11 , compared with the value of 12 shown by the Cd atom in $\mathrm{GdCd}_{3}$.

Table 4. Interatomic distances for $\mathrm{ErCd}_{3}$ and $\mathrm{YCd}_{3}$

| $\mathrm{ErCd}_{3}$ |  | $\mathrm{YCd}_{3}$ |  |
| :---: | :---: | :---: | :---: |
| $\mathrm{Er}-2 \mathrm{Er}$ | $3 \cdot 705 \AA$ | Y-2 Y | 3.76 A |
| Er-2 Cd(1) | 3.577 | Y-2 Cd(1) | $3 \cdot 56$ |
| $\mathrm{Er}-2 \mathrm{Cd}(1)$ | $3 \cdot 164$ | Y-2 Cd(1) | $3 \cdot 18$ |
| Er-2 Cd(2) | $3 \cdot 283$ | Y-2 Cd(2) | $3 \cdot 34$ |
| Er-2 Cd(2) | $3 \cdot 167$ | Y-2 Cd(2) | $3 \cdot 17$ |
| $\mathrm{Er}-4 \mathrm{Cd}(2)$ | 3.174 | $\mathrm{Y}-4 \mathrm{Cd}(2)$ | $3 \cdot 17$ |
| $\mathrm{Cd}(1)-2 \mathrm{Er}$ | $3 \cdot 577$ | Cd(1)-2 Y | $3 \cdot 56$ |
| $\mathrm{Cd}(1)-2 \mathrm{Er}$ | $3 \cdot 164$ | Cd(1)-2 Y | $3 \cdot 18$ |
| $\mathrm{Cd}(1)-2 \mathrm{Cd}(2)$ | 3.257 | $\mathrm{Cd}(1)-2 \mathrm{Cd}(2)$ | $3 \cdot 32$ |
| $\mathrm{Cd}(1)-2 \mathrm{Cd}(2)$ | $3 \cdot 190$ | $\mathrm{Cd}(1)-2 \mathrm{Cd}(2)$ | $3 \cdot 18$ |
| $\mathrm{Cd}(1)-4 \mathrm{Cd}(2)$ | $2 \cdot 964$ | $\mathrm{Cd}(1)-4 \mathrm{Cd}(2)$ | 3.00 |
| $\mathrm{Cd}(2)-1 \mathrm{Er}$ | $3 \cdot 283$ | Cd(2)-1 Y | $3 \cdot 34$ |
| $\mathrm{Cd}(2)-1 \mathrm{Er}$ | $3 \cdot 167$ | $\mathrm{Cd}(2)-1 \mathrm{Y}$ | $3 \cdot 17$ |
| $\mathrm{Cd}(2)-2 \mathrm{Er}$ | 3.174 | Cd(2)-2 Y | $3 \cdot 17$ |
| $\mathrm{Cd}(2)-1 \mathrm{Cd}(1)$ | $3 \cdot 257$ | $\mathrm{Cd}(2)-1 \mathrm{Cd}(1)$ | $3 \cdot 32$ |
| $\mathrm{Cd}(2)-1 \mathrm{Cd}(1)$ | $3 \cdot 190$ | $\mathrm{Cd}(2)-1 \mathrm{Cd}(1)$ | $3 \cdot 18$ |
| $\mathrm{Cd}(2)-2 \mathrm{Cd}(1)$ | 2.964 | $\mathrm{Cd}(2)-2 \mathrm{Cd}(1)$ | $3 \cdot 00$ |
| $\mathrm{Cd}(2)-1 \mathrm{Cd}(2)$ | 3.022 | $\mathrm{Cd}(2)-1 \mathrm{Cd}(2)$ | 3.05 |
| $\mathrm{Cd}(2)-2 \mathrm{Cd}(2)$ | 3.145 | $\mathrm{Cd}(2)-2 \mathrm{Cd}(2)$ | $3 \cdot 17$ |

The lattice constants of the orthorhombic $\mathrm{MCd}_{3}$ phases, subsequently calculated, are reported by Bruzzone et al. (1972), who noted that the cube root of the cell volume vs. the rare earth trivalent ionic radius shows a regularly decreasing linear trend from $\mathrm{TbCd}_{3}$ to $\mathrm{LuCd}_{3}$. But, if the single lattice constants of these compounds are plotted vs. the same abscissa (Fig. 4), a decreasing trend is observed for the $b$ and $c$ constants, while the $a$ constant shows an increasing trend. The inclination to form a shorter M-M distance between rare earth atoms is probably the reason for this abnormal behaviour. In fact, with decreasing rare-earth atomic dimensions, the lengthening of the elementary


Fig. 4. Lattice constants of the orthorhombic $\mathrm{MCd}_{3}$ phases (in $\AA$ ) vs. the rare earth trivalent ionic radii.
cell in the $x$ direction makes possible a large reduction in the $b$ constant, which is necessary for keeping the M-M distance short.

Apert from the stoichiometry, the atomic positions of $\mathrm{ErCd}_{3}$ are quite similar to those of NaHg .

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